Updated Smithsonian Astrophysical Observatory Ozone Monitoring Instrument (SAO OMI) formaldehyde retrieval

G. González Abad¹, X. Liu¹, K. Chance¹, H. Wang¹, T. P. Kurosu², and R. Suleiman¹

¹Harvard-Smithsonian Center for Astrophysics, Cambridge, MA, USA
²Jet Propulsion Laboratory, Pasadena, CA, USA

Correspondence to: G. González Abad (ggonzalezabad@cfa.harvard.edu)

1 November 2013 – Published in Atmos. Meas. Tech. Discuss.: 7 January 2014
Revised: 24 September 2014 – Accepted: 13 November 2014 – Published: 5 January 2015

Abstract. We present and discuss the Smithsonian Astrophysical Observatory (SAO) formaldehyde (H₂CO) retrieval algorithm for the Ozone Monitoring Instrument (OMI) which is the operational retrieval for NASA OMI H₂CO. The version of the algorithm described here includes relevant changes with respect to the operational one, including differences in the reference spectra for H₂CO, the fit of O₂–O₂ collisional complex, updates in the high-resolution solar reference spectrum, the use of a model reference sector over the remote Pacific Ocean to normalize the retrievals, an updated air mass factor (AMF) calculation scheme, and the inclusion of scattering weights and vertical H₂CO profile in the level 2 products. The setup of the retrieval is discussed in detail. We compare the results of the updated retrieval with the results from the previous SAO H₂CO retrieval. The improvement in the slant column fit increases the temporal stability of the retrieval and slightly reduces the noise. The change in the AMF calculation has increased the AMFs by 20%, mainly due to the consideration of the radiative cloud fraction. Typical values for retrieved vertical columns are between 4×10¹⁵ and 4×10¹⁶ molecules cm⁻², with typical fitting uncertainties ranging between 45 and 100%. In high-concentration regions the errors are usually reduced to 30%. The detection limit is estimated at 1×10¹⁵ molecules cm⁻².

1 Introduction

H₂CO is one of the most abundant hydrocarbons in the troposphere and a key constituent, despite its short lifetime of 1.5 h, due to photolysis and oxidation by the hydroxyl radical (OH) (Anderson et al., 1996; Brune et al., 1999). Background levels of H₂CO (concentrations below 1 ppbv) are sustained by the oxidation of methane, but over continental regions the oxidation of short-lived, non-methane volatile organic compounds (NMVOCs) from anthropogenic, biogenic, and pyrogenic origins and direct emissions from fires and industrial activities can produce enhanced concentrations of H₂CO within the boundary layer of over 5 ppbv (Houweling et al., 1998; Kanakidou et al., 2005; Lowe and Schmidt, 1983; Parrish et al., 2012). Given its short lifetime, measurements of H₂CO from satellites can be used as a proxy for NMVOC emissions and can be used to estimate top-down emission inventories (Barkley et al., 2008; Fu et al., 2007; Marais et al., 2012; Stavrakou et al., 2009a, b).

Ground-based measurements of H₂CO total columns in the ultraviolet (UV) and infrared spectral ranges are well established in the scientific community. For example, over Lauder station in New Zealand, the mean column reported by Jones et al. (2009) using high-resolution Fourier-transform infrared spectroscopy is 4.9×10¹⁵ molecules cm⁻². The same technique was used by Vigouroux et al. (2009) to measure H₂CO over Réunion Island in the Indian Ocean. They also used the multi-axis differential optical absorption spectroscopy (MAX-DOAS) technique, obtaining good agreement between the two techniques with concentrations around 4.0×10¹⁵ molecules cm⁻². MAX-DOAS H₂CO measurements (Heckel et al., 2005; Peters et al., 2012; Wagner et al., 2011; Pinardi et al., 2013) will be of great value in validating our H₂CO measurements.

H₂CO has been measured from space using both UV and the infrared radiation. Infrared solar occultation measurements from the Atmospheric Chemistry Experiment–Fourier Transform Spectrometer (ACE-FTS) (Dufour et al.,...
and emission measurements from the Michelson Interferometer for Passive Atmospheric Sounding (MIPAS) (Steck et al., 2008) are able to measure H$_2$CO profiles from the stratosphere to the upper troposphere. UV nadir measurements from the Global Ozone Monitoring Experiment (GOME), GOME-2, the Scanning Imaging Absorption spectrometer for Atmospheric CHartography (SCIAMACHY), and Ozone Monitoring Instrument (OMI) have all been used to retrieve H$_2$CO vertical columns from space (Chance et al., 2000; Palmer et al., 2001; Kurosu et al., 2004; Wittrock et al., 2006; De Smedt et al., 2008, 2012; Vrekoussis et al., 2010).

We consider the update of the operational Smithsonian Astrophysical Observatory (SAO) H$_2$CO retrieval necessary. Since 2009 there has been a significant increase of the retrieved background H$_2$CO values – e.g., Kim et al. (2011) – with data becoming noisier with time. There are new reference spectra available that need to be updated. We are also including for the first time scattering weight information in the level 2 files. We describe the theoretical basis of the SAO OMI H$_2$CO operational retrieval algorithm and the details of the setup we are using. This algorithm is used at SAO, with the necessary adjustments, for the retrieval of other trace gases (BrO (Raid Suleiman, personal communication), C$_2$H$_2$O$_2$ (Miller et al., 2014), H$_2$O (Wang et al., 2014)) from OMI.

This paper first focuses on explaining the column retrievals by describing the OMI instrument characteristics, the spectral fitting, the conversion of slant column density (SCD) to vertical column density (VCD), and the reference sector normalization. It then compares previous SAO H$_2$CO retrievals with these new ones and, finally, summarizes this study.

2 SAO OMI H$_2$CO observations

2.1 The OMI instrument

OMI (Levélé et al., 2006) was launched on 15 July 2004 onboard the NASA Earth Observing System (EOS) Aura platform into a Sun-synchronous polar orbit with an average altitude of 705 km and a period of 99 min, giving 14.7 orbits per day. The Equator-crossing time is 13:42 in the ascending node. OMI measures UV/visible (UV/Vis) solar backscattered radiation using two-dimensional charge-coupled devices (CCDs) from 270 to 500 nm in three channels: UV-1 from 264 to 311 nm, UV-2 from 307 to 383 nm, and Vis from 349 to 504 nm, with spectral resolutions of 0.63, 0.42, and 0.63 nm, respectively. We use UV-2 for H$_2$CO retrievals. In nominal operation mode, the information in the across-track direction is binned by a factor of 8, resulting in 60 binned ground pixels in the UV-2 channel. OMI has a cross-track field of view of 115$^\circ$, resulting in a swath of 2600 km, with pixel size between 13 $\times$ 24 km$^2$ at nadir and 26 $\times$ 135 km$^2$ at the swath edges, providing daily global coverage.

2.2 SAO OMI H$_2$CO retrieval

We retrieve VCDs of H$_2$CO using a two-step approach. First, we spectrally fit OMI radiances to retrieve SCDs using the basic optical differential spectroscopy (BOAS) method (Chance, 1998). Second, we convert the SCDs to VCDs using air mass factors (AMFs) calculated offline with a radiative transfer model. Finally, we apply post-processing corrections based on a reference sector over the remote Pacific Ocean (Khokhar et al., 2005).

2.2.1 Spectral fitting

SCDs of H$_2$CO are determined by direct fitting of OMI radiances (Chance, 1998). First, irradiances and radiance wavelengths are calibrated using cross correlation with a high-resolution solar spectrum (Caspar and Chance, 1997). We use the high-resolution solar spectrum recently published by Chance and Kurucz (2010). Two variables may be fitted in the wavelength calibration: spectral shift and squeeze. As they are highly correlated except for very wide fitting regions, we only fit the spectral shift. Wavelength-calibrated radiances are fitted to model spectra. The construction of the model spectra starts with the radiances from the remote Pacific Ocean, which substitute for the daily solar irradiance, scaled by a factor plus the contributions from the molecular ring correction (Chance and Spurr, 1997) and a correction accounting for the aliasing introduced from spectral undersampling computed online (Chance, 1998; Chance et al., 2005). To this term we apply the Beer–Lambert law to account for the molecular absorption from O$_3$ (Malicet et al., 1995), NO$_2$ (Vandaele et al., 1998), BrO (Wilmouth et al., 1999), the O$_2$–O$_3$ collision complex (Thalman and Volkamer, 2013), and H$_2$CO (Chance and Orphal, 2011), all convolved with the OMI slit function measured prior to launch (Dirksen et al., 2006). After accounting for the molecular absorption, we include the common mode spectrum obtained by averaging the fitting residuals for each orbit. We finish the construction of the model spectra by considering two low-order closure polynomials that help to account for low-frequency features of the measured radiance in the fitting window (Eq. 1).

$$I = \left[ (a I_0 + \sum_i a_i X_i) e^{-\sum_j a_j X_j} + \sum_k a_k X_k \right] \sum_n a_n X_n + \sum_m a_m X_m$$

(1)

In Eq. (1), $a$ is the factor applied to $I_0$, the solar irradiance, or in this case the radiance reference from the remote Pacific Ocean; $a_i X_i$ are terms added before the trace gas contributions, including the ring spectrum and the undersampling correction; $a_j X_j$ are the Beer–Lambert law contributions from the trace gases; $a_k X_k$ are the terms added afterwards, which in this particular case is the common mode spectrum; and $a_n X_n$ and $a_m X_m$ are the scaling and baseline.
closure polynomials. The inclusion of the common mode has the intention of reducing the fitting residuals without affecting the values of the retrieved slant columns since its correlation with the H$_2$CO cross sections is small. Figure 1 shows the scatterplots of the SCDs and the fitting uncertainties for Orbit 11118 comparing the retrieval when we use or do not use the common mode. The correlation values between the two setups is high (0.9996 for the SCDs and 0.9987 for the fitting uncertainties), with small changes in the SCD values. As said before, the main reason to include the common mode is to reduce the root mean square (rms) and therefore the fitting uncertainties, which, as shown in the right panel of Fig. 1, are reduced on average 2%.

Figure 2 shows the correlations of the other molecules included in the fitting with H$_2$CO for Orbit 8537. We are also showing the correlation with the common mode. Correlations of H$_2$CO with O$_3$ (at 228 and 295 K), with NO$_2$, O$_2$–O$_2$ collision complex, and the common mode spectrum are all small. The orbit-averaged correlations are 0.15, 0.15, 0.13, 0.05, and 0.01, respectively. The common mode is weakly correlated with H$_2$CO. As expected from the correlation of cross sections between BrO and H$_2$CO, the correlation of the retrieved column amount between them is strong. The average correlation coefficient for the whole orbit is 0.44. As a result, H$_2$CO retrieval is affected if BrO is present in significant amounts in the atmosphere.

Instead of using the traditional solar irradiance as $I_0$ in Eq. (1), we use radiances over the remote Pacific Ocean obtained from an orbit as close to 165° W as possible, measured within 24 h of the orbit to be processed. We build the radiance reference by averaging the radiance spectra over the Pacific between 30° S and 30° N. Then we use this radiance reference as $I_0$. Due to the weak absorption of H$_2$CO, the use of the radiance reference over the Pacific helps to address some instrument issues observed in previous versions of the H$_2$CO retrieval, including the well-known striping problems of OMI (Veihelmann and Kleipool, 2006). The characteristics of the spectral fit, reference spectra, closure polynomials, and online computed corrections are summarized in Table 1.

![Figure 1. Fitting results of H$_2$CO SCDs for OMI Orbit 11118 with and without common mode (17 August 2006 overpass over South America).](image1)

![Figure 2. Correlations of H$_2$CO with the rest of molecules considered in the spectral fitting and the common mode spectrum for Orbit 8537 (21 February 2006 overpass over Africa and Europe).](image2)
Table 1. Fitting window and parameters for deriving H$_2$CO slant column density.

<table>
<thead>
<tr>
<th>Parameter</th>
<th>Details</th>
</tr>
</thead>
<tbody>
<tr>
<td>Fitting window</td>
<td>328.5–356.5 nm</td>
</tr>
<tr>
<td>Radiance reference spectrum</td>
<td>Computed online over the remote Pacific Ocean between 30$^\circ$ N and 30$^\circ$ S</td>
</tr>
<tr>
<td>Baseline polynomial</td>
<td>3rd order</td>
</tr>
<tr>
<td>Scaling polynomial</td>
<td>3rd order</td>
</tr>
<tr>
<td>Instrument slit function</td>
<td>Hyper-parameterization of pre-flight measurements</td>
</tr>
<tr>
<td>H$_2$CO cross sections</td>
<td>Chance and Orphal (2011), 300 K</td>
</tr>
<tr>
<td>O$_3$ cross sections</td>
<td>Malicet et al. (1995), 228 and 295 K</td>
</tr>
<tr>
<td>NO$_2$ cross sections</td>
<td>Vandeale et al. (1998), 220 K</td>
</tr>
<tr>
<td>BrO cross sections</td>
<td>Wilmouth et al. (1999), 228 K</td>
</tr>
<tr>
<td>$O_2$–$O_2$ collision complex cross sections</td>
<td>Thalman and Volkamer (2013), 293 K</td>
</tr>
<tr>
<td>Molecular ring cross sections</td>
<td>Chance and Spurr (1997)</td>
</tr>
<tr>
<td>Undersampling correction</td>
<td>Computed online Chance et al. (2005)</td>
</tr>
<tr>
<td>Residual common mode spectrum</td>
<td>Computed online between 30$^\circ$ N and 30$^\circ$ S</td>
</tr>
</tbody>
</table>

Figure 3. Fitting of H$_2$CO in OMI Orbit 11118 (17 August 2006 overpass over South America) for high (top), medium (middle), and low (bottom) amounts. The black line is the fitted H$_2$CO optical depth, and the red line is the fitted H$_2$CO optical depth plus the fitting residuals.

is presented here. Figure 3 shows the residual and the results of the fitting for three pixels of Orbit 11118 (17 August 2006) over South America. The thick black line is the fitted H$_2$CO optical depth, and the thin red line is the fitted H$_2$CO optical depth plus the residual, shown here to provide an idea of the fitting quality.

Typical values for the retrieved H$_2$CO SCDs are in the range from $4 \times 10^{15}$ to $6 \times 10^{16}$ molecules cm$^{-2}$. Fitting uncertainties for individual pixels range from $6 \times 10^{15}$ molecules cm$^{-2}$ ($\sim$ 30%) for pixels with high concentrations (above $2 \times 10^{16}$ molecules cm$^{-2}$) to 100% or larger for pixels with low concentrations. These fitting uncertainties are calculated for the variable $j$ considering the diagonal term of the covariance matrix for the fitting variable $C_{j,j}$ and the fitting residual weighted by the number of degrees of freedom (Eq. 2). Other sources of error in the fitting of the slant columns such as forward model errors or errors associated with parameters not optimized during the fitting are not considered in the calculation of the fitting uncertainties. Fitting rms values vary between $0.4 \times 10^{-3}$ and $2.0 \times 10^{-3}$, corresponding to detection limits between $\sim 6 \times 10^{15}$ and $\sim 3 \times 10^{16}$ molecules cm$^{-2}$.

$$\delta X_j^2 = \text{RMS}^2 \left( \frac{m}{m-n} \right) (C_{(j,j)} C_{(j,j)})$$

(2)

2.2.2 Vertical column density determination

The second step of our retrieval converts SCDs into VCDs. The photons reaching OMI are either reflected from the Earth’s surface and transmitted or scattered to OMI or directly scattered to OMI without reaching the Earth’s surface. The retrieved SCDs depend on the viewing geometry for each pixel (solar zenith angle, viewing zenith angle, and relative azimuth angle), the surface albedo, and the state of the atmosphere (clouds, trace gases concentrations, temperature, air density, pressure, etc.). The SCDs are related to the VCDs by AMFs. The relationship between SCDs and VCDs is

$$\text{VCD} = \text{SCD}/\text{AMF}.$$  

(3)

To determine AMFs, we follow the approach described in Palmer et al. (2001), summarized by Eq. (3).

$$\text{AMF} = \int w(z, \theta_0, \theta, \varphi, \alpha_s, h_s, C_f, C_h) S(z, \text{lat}, \text{lon}, \text{month}) dz$$  

(4)

In this formulation, the scattering weights ($w$’s) account for the sensitivity of the satellite measurements to the presence
of H₂CO molecules at different altitudes. They are functions of altitude \( z \), solar zenith angle \( \theta_0 \), viewing zenith angle \( \theta \), relative azimuth angle \( \phi \), surface albedo \( a_s \), terrain height \( h_s \), effective cloud fraction \( C_T \), and cloud-top altitude \( C_h \). For an optically thin absorber such as H₂CO, the \( w \)'s are independent of the column abundance of the species. The shape factor \( S(z) \), dependent on the location and time of the measurement, represents the normalized vertical profile of the absorbing molecule and provides a priori estimation of the vertical distribution of H₂CO. The shape factor for a particular location and time is defined in Eq. (4).

\[
S(z) = \frac{x_a(z)}{\int_{\text{atm}} x_a(z) dz}
\]

(5)

\( x_a \) is the a priori vertical profile of H₂CO partial columns. We have calculated vertical profiles using the GEOS-Chem (Bey et al., 2001) chemical transport model (Version v9-01-03), which is driven by Goddard Earth Observing System-5 (GEOS-5)-assimilated meteorological data products interpolated onto 47 vertical levels at a spatial resolution of 2° latitude × 2.5° longitude. We have developed a monthly mean climatology using GEOS-Chem simulations between 13:00 and 14:00 local time (LT) in 2007. Figure 4 shows a set of vertical profiles used in the AMF calculation for Orbit 11118 at different latitudes.

We have calculated the scattering weights using the radiative transfer model VLIDORT version 2.4RT (Spurr, 2006). \( w(z) \) usually depends on the wavelength and the viewing geometry. However, for solar zenith angles below 70°, the ozone absorption is not too strong in the 328.5–356.5 nm fitting window, the variation of \( w(z) \) with wavelength is small (less than 5%). Therefore, we have used the scattering weights at only one optimized wavelength, 340 nm. For solar zenith angles above 70°, the ozone absorption becomes large enough to make the assumption of an optically thin atmosphere not valid (we concentrate on SZA < 70°). We have prepared a set of weighting function look-up tables that are parameterized by the surface albedo and the relative azimuth angle, with entries for solar zenith angles (12), viewing zenith angles (8), altitudes (48), and cloud-top/surface pressures (6). These parameters are summarized in Table 2. OMCLD02, the OMI cloud product we have used to obtain the cloud information (Acarreta et al., 2004), assumes the clouds to be opaque Lambertian surfaces at a single cloud height \( C_h \) with a cloud albedo \( a_c \) of 0.8–1.0. At the retrieval partly cloudy pixels are assumed to be a mixture of clear-sky and cloudy scene (\( a_c = 0.8 \)) with an effective cloud fraction \( C_T \) calculated using the independent pixel approximation (Martin et al., 2002). Under this approximation, the scattering weight for an inhomogeneous pixel with a given viewing geometry becomes

\[
w = (1 - \Phi) \cdot w_{\text{clear}}(a_s, h_s) + \Phi \cdot w_{\text{cloud}}(a_c, C_h),
\]

(6)

where \( w_{\text{clear}} \) is the scattering weight for the clear-sky part of the pixel, \( a_s \) is the surface albedo, \( h_s \) is the terrain height, \( w_{\text{cloud}} \) is the scattering weight of the cloudy part of the pixel, \( C_h \) is the cloud height, and \( \Phi \) is the intensity-weighted cloud fraction or radiative cloud fraction calculated as

\[
\Phi = \frac{C_T \cdot I_{\text{cloud}}}{(1 - C_T) \cdot I_{\text{clear}} + C_T \cdot I_{\text{cloud}}}
\]

\( I_{\text{clear}} \) and \( I_{\text{cloud}} \) are the radiance intensities for clear-sky and cloudy scenes at the top of the atmosphere. The values of \( I_{\text{clear}} \) and \( I_{\text{cloud}} \) are calculated using VLIDORT and stored in a look-up table with the same parameter grid as the weighting function look-up table. The \( a_s \) used to evaluate \( w_{\text{clear}} \) and \( I_{\text{clear}} \) is obtained from the most frequent Lambertian equivalent reflectance (LER) field in version 3 of the OMI surface reflectance climatology (Kleipool et al., 2008). We calculate \( a_s \) by interpolating the values contained in the climatology to the wavelength where the AMFs are calculated, 340 nm. For the terrain height, \( h_s \), we use the information provided with the OMI level 1B product (OML1BRUG). The description of clouds is consistent with the assumptions used in the cloud retrieval algorithm (Acarreta et al., 2004; Stammes et al., 2008) from which we obtain \( C_T \) and \( C_h \) distributed in the Aura OMI Cloud Data Product-OMCLD02. Figure 5 shows the scattering weights for three pixels of Orbit 11118 with low, medium, and high effective cloud fractions. Once the scattering weights are calculated, the averaging kernels (AKs) (Rodgers and Connor, 2003) are straightforward to be approximated using the formulation shown below for DOAS retrievals (Eskes and Boersma, 2003):

\[
\text{AK}(z) = \frac{w(z)}{\text{AMF}}.
\]

w is calculated as in Eq. (5), and AMFs are described by Eq. (3). Model comparisons with H₂CO VCDs, assimilation of OMI H₂CO products, and validation with H₂CO profiles
Table 2. Parameters of the weighting function look-up table.

<table>
<thead>
<tr>
<th>Parameter</th>
<th>Grid values</th>
</tr>
</thead>
<tbody>
<tr>
<td>Solar zenith angle</td>
<td>0, 15, 30, 45, 60, 70, 77, 81, 84, 86, 88, 89 [°]</td>
</tr>
<tr>
<td>Viewing zenith angle</td>
<td>0, 15, 30, 45, 60, 70, 75, 80 [°]</td>
</tr>
<tr>
<td>Altitude</td>
<td>0.06, 0.19, 0.32, 0.45, 0.59, 0.73, 0.86, 1.00, 1.15, 1.29, 1.44, 1.58, 1.76, 1.99, 2.25, 2.52, 2.79, 3.08, 3.44, 3.90, 4.38, 4.88, 5.42, 5.98, 6.59, 7.24, 7.95, 8.86, 9.96, 11.04, 12.11, 13.15, 14.19, 15.22, 16.24, 17.26, 18.81, 20.93, 23.12, 25.42, 29.16, 34.65, 40.93, 48.04, 55.86, 64.16, 74.49, 80.79 [km]</td>
</tr>
<tr>
<td>Cloud-top/surface pressure</td>
<td>1.00, 0.83, 0.63, 0.52, 0.37, 0.24 [atm]</td>
</tr>
</tbody>
</table>

Figure 5. Example of calculated scattering weights for three pixels with different cloud fractions corresponding to the profiles plotted in Fig. 3. For cloud fractions equal to 1 the scattering weights below the cloud-top pressure become 0 (green line), while for other values there is a contribution below the cloud in the air mass factor calculation (red and black lines) proportional to the radiative cloud fraction.

should make use of the scattering weights and the used H$_2$CO vertical profiles, so we provide them along with the columns.

We have performed a sensitivity analysis to assess the behavior of the AMFs calculations with respect to the H$_2$CO vertical profile, the wavelength, the surface albedo, the viewing geometry, and the cloud parameters. We have defined a reference scenario for the AMF calculation considering a cloud-free pixel, with a surface albedo of 0.02, a solar zenith angle of 30°, and a viewing zenith angle of 0°. The wavelength considered for the standard calculation is 340 nm.

As mentioned above, the variation of the AMFs in function of the wavelength in the range of the fitting window is small, between 1 and 5%. It is true and worth mentioning that at shorter wavelengths, below 330 nm, the AMFs have a stronger dependency on the wavelength. This is one of the reasons to keep the lower limit of the fitting window on 328.5 nm.

In the UV spectral range, most surfaces reflect between 1 and 10% of the incoming light. Kleipool et al. (2008) estimated the errors for the LER in the UV to be around 40%. We can then estimate the errors we are carrying over from the albedo data into the AMF calculation for different surfaces. For example, the AMF of a typical surface with an albedo of 0.03, such as a forest, will have a 6% error associated with the OMLER uncertainty. But for brighter surfaces, the error will be around 14% for the Sahara desert with an albedo of 0.1 and around 26% for a snow surface with an albedo of 0.8. Therefore, we recommend the users to exercise caution when using retrievals over high albedo surfaces.

As mentioned before the dependence of the AMF in function of the SZA and VZA is strong. For example the AMF for our reference scenario can vary between 2 and 6. Fortunately the geometries of the satellite observations are determined precisely and the errors associated to the inaccuracy of the viewing geometry are neglected.

The cloud pressure has an important impact in the value of the computed AMFs. Considering the reference scenario and a pixel fully covered with clouds the AMFs can range from 0.6 to 1.2 when the cloud pressure increases from 800 to 900 hPa. The cloud product we use is derived using the O$_2$–O$_2$ absorption band at 477 nm (Acarreta et al., 2004). An alternative to this product is the OMCLDRR cloud product, derived from rotational Raman scattering around 350 nm (Joiner and Vassilkov, 2006). The difference between these two data sets was estimated to be between 2 and 45 hPa (Sneep et al., 2008). After propagating this difference in the cloud pressure to the calculations of the AMFs the values of the AMFs change around 20%. The other cloud parameter we need to consider is the cloud fraction. Acarreta et al. (2004) estimated the systematic uncertainty of the cloud fraction to be ±0.1. In the case of considering a cloud pressure of 800 hPa a change in the cloud fraction of ±0.1 implies changes in the AMFs of around 25%.
We have also tested the impact of the a priori profiles of H$_2$CO. Assuming a bias in the GEOS-Chem profiles of ±10 % (Millet et al., 2006; Barkley et al., 2011) we have found that the AMF may change for different regions up to 16%.

To summarize, the errors associated with AMF calculations are estimated to be around 35 %, dominated by the cloud parameters, the surface albedo, and the H$_2$CO vertical profiles. The presence of aerosols in the atmosphere is an important factor, but we have not explicitly included them in our look-up tables to keep them consistent with the OMCLDO2 cloud product (Acarreta et al., 2004). The AMF uncertainties combined with the estimated errors on the SCDs (30–100 %) result in uncertainties associated with the VCDs between 45 and 105 %.

2.2.3 Post-processing normalization of retrieved columns

We derive a daily post-processing normalization for the retrieved columns according to the concentrations simulated by GEOS-Chem (Bey et al., 2001) over the remote Pacific, following an approach similar to Khokhar et al. (2005) and De Smedt et al. (2008). In this region, the H$_2$CO concentrations are assumed to be at background levels, their only source being CH$_4$ oxidation.

We define the reference sector as the region in-between 140 and 160° W, and 90° N and 90° S. Using 2.0° latitude by 2.5° longitude model results in 2007, we have built a monthly mean climatology of VCDs by longitudinally averaging concentrations for a given latitude and then latitudinally interpolating the results to a grid of 500 points extending from 90° N to 90° S with a resolution of 0.36°. With these climatological values, taking into consideration the viewing geometry and atmospheric conditions of each pixel (i.e., AMF), we correct the effects of the row anomaly and stripping and normalize the retrieved values to the reference sector to remove the artificially increasing background H$_2$CO values using the following approach.

We calculate a pixel-by-pixel correction (Eq. 8) over the remote Pacific Ocean (radiance reference granule) for the pixels whose SCD fitting has converged successfully above the noise level, and the retrieved SCD plus 2 times the fitting uncertainty are above 0. These pixels are identified in the file product as good pixels by having a main data quality flag of 0. Other possible values for the main quality flag are 1 for suspect pixels – with negative values for the retrieved column plus 2 times the fitting uncertainty but positive values for the column plus 3 times the fitting uncertainty, or with convergence of the fitting at the noise level – and value 2 for bad pixels, which are pixels where the fitting has not converged, or the retrieved values plus 3 times the fitting uncertainty remain below 0.

Correction($i$, $j$) = OMI$_{SCD,Pacific}(i, j)$

− (GEOS$_{Chem, ReferenceSector}$(lat) × AMF($i$, $j$))

(9)

$i$ is the a cross-track position (longitude) within an OMI orbit, $j$ is the line position (latitude), and lat is the latitude at the center of pixel ($i$, $j$) of the orbit.

After working the pixel-by-pixel correction using the radiance reference orbit, we assume that the bias between the retrieved columns and the modeled columns over the remote Pacific Ocean (Correction($i$, $j$)) is representative of a longitudinally invariant bias that affects the retrievals everywhere else on the globe. Therefore, we calculate a latitude-dependent correction of this bias for each cross-track position using the median values of the corrections calculated in Eq. (8) falling within two points of the reference sector latitude grid. Finally we apply this bias to the rest of the orbits by interpolating the values of the median bias to the corresponding latitude at the center of each pixel.

We attribute this bias mainly to instrumental issues whose effects are weaker than the row anomaly, and to correlations between BrO and H$_2$CO absorption cross sections, which are especially significant at high latitudes. One latitude-dependent bias correction is calculated for each radiance reference orbit. In fact this implies that the bias correction is calculated daily. We consider this bias correction to be orbit invariant for all the retrievals performed with the same radiance reference (within 24 h difference). Therefore we apply it to all OMI columns retrieved from different granules over the globe (typically between 7 and 17 granules using a particular radiance reference). The final corrected columns are calculated as

\[ \text{VCD}_{\text{corrected}}(i, j) = \frac{\text{SCD}(i, j) - \text{Correction}(i, j)}{\text{AMF}(i, j)} \]

(10)

Figure 6 illustrates the GEOS-Chem results used in the calculation of the background levels for August 2007. Figure 7 shows the effect of the reference sector correction for the reference sector (140 to 160° W and 90° N to 90° S) for each month. We can see how at high latitudes in the winter months we consistently retrieve high values of H$_2$CO when compared with the GEOS-Chem reference sector slant columns (left panels). It is therefore no surprise that the reference sector correction is strongest for the winter months in the high latitudes. We consider the month-to-month variation of the GEOS-Chem climatology to be representative of the changes in the background values, in comparison to the rapid spatial and temporal variation in concentrations over the continental hot spots.

Even though the correction helps us to remove unrealistic biases associated with the instrument and correlations between the H$_2$CO and BrO cross sections, it introduces external a priori information subject to uncertainties and assumes that H$_2$CO over the reference sector does not change from
3 Comparison between previous and current H$_2$CO SAO product

We have changed the spectral fitting for the new version of the retrieval by updating some cross sections to newly available data, by considering the O$_2$–O$_2$ collisional complex, and by changing the fitting window. Table 3 shows the mean SCD, mean fitting uncertainty, mean rms, and the percent of pixels flagged as good (main quality flag equal to 0) for OMI Orbit 11148 using different retrieval configurations. First we report the result from the old SAO retrieval and then the result after changing each one of the items listed below. The last line of Table 3 shows the results for the new retrieval. The mean SCD has changed from 7.52 to 11.02 × 10$^{15}$ molecules cm$^{-2}$, the fitting uncertainty from 11 to 8.87 × 10$^{15}$ molecules cm$^{-2}$, the rms from 1.31 to 1.22 × 10$^{-3}$, and the number of good pixels rises from 98.95 to 99.64 %.

The list of changes in the spectral fitting includes the high-resolution solar spectrum from Chance and Spurr (1997) to Chance and Kurucz (2010), the H$_2$CO cross sections from Cantrell et al. (1990) to Chance and Orphal (2011), and the use of two O$_3$ cross sections at 228 and 295 K instead of only one at 228 K (Malicet et al., 1995). Another change is that we are fitting the O$_2$–O$_2$ collisional complex cross section (Thalman and Volkamer, 2013), while in the old SAO retrieval the O$_2$–O$_2$ collisional complex was not considered. The O$_2$–O$_2$ collisional complex has two significative features at 342 and 360 nm that are relevant in the H$_2$CO fitting window. For that reason we consider it necessary to include it in the fitting. We have also modified the fitting window from 327.5–356.5 nm in the old SAO retrieval to 328.5–356.5 nm in the new retrieval.

To compare the long-term series of both retrievals, 3 years of data are used; years 2006 and 2007, near the beginning of the mission, are representative of a time when the previous SAO product is quite stable, while year 2012 is representative of a time when the previous SAO product is affected by problems, including the row anomaly and increasing background concentrations.

In general, the new product is less noisy, with smaller fitting uncertainties. This can be seen for the particular case of Orbit 43214 (August 2012) in Fig. 8. The decrease in uncertainty is evident with almost all pixels having an absolute uncertainty smaller than that in the older version of the SAO retrieval. There is also significant reduction of unrealistically high SCDs.

The general increase in the retrieved columns in the previous SAO retrieval is so large that it is difficult to use it after 2010 (Kim et al., 2011). In the new retrieval, the situation is greatly improved. In Fig. 9 (for Orbit 43214) we have plotted the retrieved VCDs for the old retrieval, the new one, and the new one with the reference sector correction applied to show how the situation has improved. The first panel, top left, gives VCDs retrieved with the old algorithm. Extremely high values (over 2.0 × 10$^{16}$) are common over the continents and
Table 3. Spectral fitting results for Orbit 11118. Results from the old SAO retrieval and from different new retrieval settings are included to illustrate the change induced by each one of them.

<table>
<thead>
<tr>
<th></th>
<th>Mean SCD molecules cm⁻²</th>
<th>Mean uncertainty molecules cm⁻²</th>
<th>Mean rms</th>
<th>% good pixels</th>
</tr>
</thead>
<tbody>
<tr>
<td>Old SAO</td>
<td>7.52 × 10^{15}</td>
<td>1.10 × 10^{16}</td>
<td>1.31 × 10^{-3}</td>
<td>98.95</td>
</tr>
<tr>
<td>New fitting window</td>
<td>8.41 × 10^{15}</td>
<td>1.03 × 10^{16}</td>
<td>1.31 × 10^{-3}</td>
<td>99.14</td>
</tr>
<tr>
<td>New solar reference</td>
<td>5.56 × 10^{15}</td>
<td>1.02 × 10^{16}</td>
<td>1.34 × 10^{-3}</td>
<td>98.53</td>
</tr>
<tr>
<td>New H₂CO cross sections</td>
<td>4.48 × 10^{15}</td>
<td>9.14 × 10^{15}</td>
<td>1.33 × 10^{-3}</td>
<td>98.31</td>
</tr>
<tr>
<td>New O₃ cross sections</td>
<td>4.95 × 10^{15}</td>
<td>9.60 × 10^{15}</td>
<td>1.26 × 10^{-3}</td>
<td>97.67</td>
</tr>
<tr>
<td>O₂–O₂ cross sections</td>
<td>7.03 × 10^{15}</td>
<td>9.85 × 10^{15}</td>
<td>1.29 × 10^{-3}</td>
<td>98.77</td>
</tr>
<tr>
<td>New SAO</td>
<td>1.02 × 10^{16}</td>
<td>8.87 × 10^{15}</td>
<td>1.22 × 10^{-3}</td>
<td>99.64</td>
</tr>
</tbody>
</table>

Figure 8. Scatterplots for Orbit 43214 comparing the old SAO retrieval and new SAO retrieval of SCDs and SCD fitting uncertainties.

Equatorial ocean areas, while there are areas of negative values at high latitudes.

With the new algorithm without the reference sector correction, top right panel, it is evident that there is a significant reduction in the concentrations almost everywhere. The contrast between high- and low-concentration regions increases. The new retrieval also has fewer pixels with negative values at high latitudes. After we have applied the reference sector correction (bottom left panel), the situation further improves. As mentioned earlier, the new retrieval still shows large negative concentrations – which, however, are smaller in the new version – at high latitudes due to correlations with BrO. Applying the reference sector correction helps to minimize this effect. Finally, we plot the difference between the VCDs retrieved with the old and the new products in the bottom right panel. The concentrations are reduced almost everywhere in the orbit (blue colors). We only see increases in concentrations where the old retrieval is consistently negative. There is also an improvement for pixels affected by the row anomaly in the new retrieval. However, we do not recommend using pixels affected by the row anomaly.

We now compare level 3 monthly averages to better assess changes in the concentrations and trends of the new retrieval with respect to the old SAO retrieval. To compute monthly averages, we only considered pixels with cloud fractions below 40% and unaffected by the row anomaly. The level 3 data are calculated on a 0.2° × 0.2° grid where values of the OMI pixel are weighted by the area falling within the grid box and by the fitting uncertainty.

We have defined six regions in this study. Background levels over the remote Pacific Ocean (A), southeastern USA (B), the Amazon Basin (C), Europe (D), and southeast Asia (E) plus an equatorial Pacific Ocean area which includes only the latitudes between −20° S and 20° N to assess the evolution of the SCD standard deviation in a remote area. The areas used for each of these regions are shown in Fig. 6. We have focused our analysis on 2006, 2007, and 2012. While the old SAO retrieval is stable until the end of 2007, by 2012
it is clearly producing unreasonable columns over the Pacific and an overall increase of the retrieved concentrations worldwide. The selected 3 years are thus appropriate for comparing the old and the new results.

Figure 10 shows the results of the time series comparison. Overall, the new retrieval (red and blue lines in the left panels) is more stable over time than the old retrieval (black and green lines in the left panels). The VCDs plotted for the new retrieval are the ones with the reference sector correction applied to them. In comparison with the old SAO retrieval the new one does not show a large increasing trend in the SCD or VCD concentrations. The background concentrations over the Pacific Ocean remain almost the same for the new retrieval, while the old SAO retrieval exhibits an increase of $0.3 \times 10^{16}$ molecules cm$^{-2}$ between the end of 2007 and the beginning of 2012. However there is still some increase, albeit greatly reduced. We have not been able to attribute it to any specific cause after analyzing possible anisotropies introduced by the row anomaly in the viewing geometry of the pixels. What we can say is that we do not observe a degradation of the fitting residuals (orange and purple lines of the right side panels in Fig. 10) or increases in the standard deviations (black and red lines in the right side panels) in the new retrieval in opposition to what was observed in the old retrieval. The other regions have similar seasonality in the old and new retrievals, but the sharp increase in the background levels observed in the old retrieval from 2007 to 2012 is not present in the new retrieval, especially in the southeastern USA and southeast Asia. The retrieved columns in the new version do not show significant jumps between 2007 and 2012 for the rest of the regions analyzed. The main remaining issue in the new retrieval is the too-high winter concentrations over Europe. We recommend at present the values for December, January, and February over northern Europe to be used with caution.

In the left side panels of Fig. 10 we have also included the AMF for the old and new retrieval (orange and purple lines). As mentioned above, the use of the radiative cloud fraction instead of the effective cloud fraction translates to a general increase of the AMFs of around 20%.

Figure 11 shows the results for the new retrieval without the reference sector correction applied in the middle panels,
with the reference sector correction applied in the right panels, and the old SAO retrieval in the left panels. These plots are seasonal averages for DJF, MAM, JJA, and SON for the years 2006 and 2012. The reduction in the background values is evident in the new retrieval. Hot spots are more spatially resolved in the new retrieval, e.g., the hot spot over Central America in MAM. The high concentrations over the Mediterranean Sea in JJA are removed. To illustrate the reduction of the increasing trend in the new retrieval, Fig. 12 shows the difference of seasonally averaged slant columns between 2006 and 2012 (2012–2006) for the old (left column) and new (right column) retrievals. Apparently, the new retrieval has reduced the increasing trend significantly. The differences are well over 100 % in the old version for most continental and oceanic regions in all seasons. While there is still an increasing trend in the new retrieval, it has been decreased to between 0 and 50 % in most regions and seasons.

4 Conclusions

We have presented and discussed the new SAO OMI H$_2$CO retrieval, paying especially attention to the changes with respect to the old operational algorithm. We have updated the H$_2$CO cross sections and the high-resolution solar spectra. We have also included a second O$_3$ cross sections at 295 K and the O$_2$–O$_2$ collisional complex. We have implemented an iterative fitting residual analysis to remove spectral pixels with high fitting residuals. We have slightly modified the fitting window and used a reference sector over the remote Pacific Ocean to correct for possible biases and their temporal drift in the retrieved slant columns. We have also updated the AMF calculation to include the radiative cloud fraction. This change has increased the value of the AMFs by 20 % on average.

The columns retrieved by the new SAO OMI H$_2$CO algorithm show more stability over time. The increasing trend in concentrations is greatly reduced, as is the increase in the rms and therefore the fitting uncertainties and the standard deviations of the columns. This new product greatly extends the volume of useful H$_2$CO data from OMI beyond 2008 and also provides new data to facilitate its use. The scattering weights used in the SCD-to-VCD inversion and the associated H$_2$CO vertical profiles are now included in the level 2 data files for use in model comparisons. We now include a quality flag in our product for problems caused by instrument row anomaly. As a general rule, we recommend users not to use pixels affected by the row anomaly.

The new operational algorithm will be version 3.2.0.3 in the algorithm development sequence. Generically we will refer to it as version 3. A detailed error analysis is in preparation. We are also working on further validation of our new product.

Acknowledgements. This study is supported by NASA Atmospheric Composition Program/Aura Science Team (NNX11AE58G) and the Smithsonian Institution. The Dutch–Finnish OMI instrument is part of the NASA EOS Aura satellite payload. The OMI project is managed by NIVR and KNMI in the Netherlands. We acknowledge the OMI International Science Team for providing OMI data used in this study.

Edited by: F. Boersma

References


Diem et al., 2005: Evaluation of the OMI H2CO retrieval


